

ITP lunch seminar

Introduction

At fixed order

$$\sigma = \int_n d\phi_n \left| A_n^{(0)} + A_n^{(1)} + A_n^{(2)} + \dots \right|^2 \quad S(p_i)$$

$$+ \int_{n+1} d\phi_{n+1} \left| A_{n+1}^{(0)} + A_{n+1}^{(1)} + \dots \right|^2 \quad S(p_i)$$

$$+ \int_{n+2} d\phi_{n+2} \left| A_{n+2}^{(0)} + \dots \right|^2 \quad S(p_i)$$

$\swarrow L_0$ $\swarrow M_0$ $\swarrow N_0$
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need to use numerical integration for phase space as measurement function S too complicated / not known
 real matrix element

$$|A_{n+1}^{(0)}|^2 \sim \left| \sqrt{\frac{1}{E_\gamma}} + \dots \right|^2 \sim \frac{1}{E_\gamma^2} \frac{1}{1 - \cos\theta}$$

$E_\gamma \rightarrow 0$
 $\rightarrow \infty$
 for massive fermions $\cos\theta \rightarrow 1$ is fine as $\beta < 1$

How to do that in d -dim?

Subtraction & slicing

Idea 1 slicing

$$\int d\sigma_{n+1} = \int_0^{\omega_c} d\sigma_{n+1} \uparrow \int_{\omega_c} d\sigma_{n+1}$$

approximate \hookrightarrow
&
analytic

\hookrightarrow exact
&
numeric

$$= \frac{1}{\epsilon} + \log \omega_c - \log \omega_c + \text{finite}$$

large cancellation as $\frac{2\omega_c}{\sqrt{s}} \approx 10^{-6}$

\leadsto end up wasting most of your MC points on something that cancels anyway

Idea 2: subtraction

$$\int d\sigma_{n+1} = \int dCT + \int (d\sigma_{n+1} - dCT)$$

analytic \hookrightarrow

\hookrightarrow locally finite
& well-behaved

(can be matched with slicing)

\leadsto numeric

main advantage of subtraction:
+ local cancellation rather than global

(2-6x faster [2003-01014]
10x faster [Dittmaier])

- generally more complicated

First scheme @ NLO FKS '95, later CS '96

NNLO Antenna '07

STRIPPER '10

Projection-to-Bar '15

residue-improved sec. decomp '17

Turin scheme '18

side note

$$\int \frac{d^d k}{(2\pi)^d} \text{Sun} + \int \frac{d^{d-1} k}{(2\pi)^{d-1}} 2k^0 dCT = \text{finite}$$

use loop-tree?

For QED dCT & $\int dCT$ is known to
all orders (FKS^e)

↪ can calculate fixed-order
cross-section & histograms to N^eLO

Event generation

experimentalists need generators, i.e.

events $\{P_1, P_2, P_3, \dots, P_n\}$, w , for detector-sim.

trivial in principle, just dump all MC points to file (garden hose approach)

Problem: weights w not uniform, many orders of magnitude & negative

$$\int d\sigma_{n+1} = \underbrace{\int dCT}_{\geq 0} + \underbrace{\int (d\sigma_{n+1} - dCT)}_{\text{whatever}}$$

slicing actually has an advantage here...

not a problem for $d\sigma/dx$ b/c S is fast $\mathcal{O}(\mu s)$

detector sim $S \mathcal{O}(\text{min-hours})$

\rightarrow minimise # of events

estimate: say we have N events $w = zc$
 $r \cdot N$ of which are negative.

variance

$$\sim \frac{\sum_i w_i^2}{(\sum_i w_i)^2} = \frac{c^2 N}{[N(1-r)c - Nr c]^2} = \frac{1}{N} \frac{1}{(1-2r)^2}$$

need $\frac{1}{(1-2r)^2}$ more points...

LHC example

pp \rightarrow $\gamma\gamma$ + jets	5x
Z + 1 jet	2500x
Z + 2 jet	6000x
Z + 3 jet	9000x
W + 5 jet	14000x

need to reduce r !

Cellular resampling

Andersen, Maier '21

Andersen, Maier, Maître '23

github.com/a-maier/cres

Two observations:

- cross sections are positive
 $\forall \text{Region } \int_{\mathcal{R}} d\sigma > 0$
- experiments have finite resolution
 $\leadsto \text{KLN}$

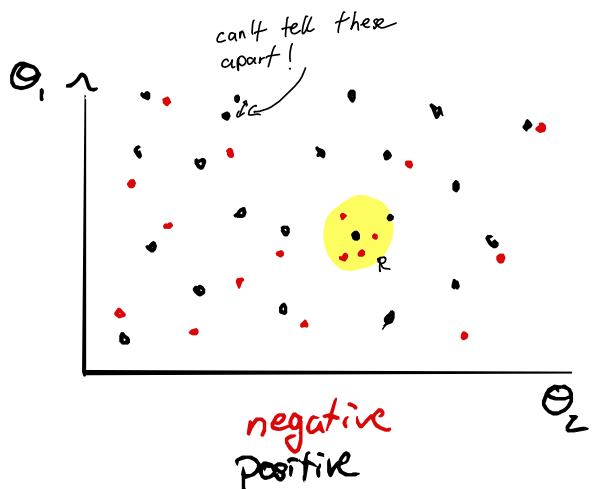
$\forall \text{ event } e_1 \exists e_2 \text{ st. } e_1 \text{ \& } e_2$
are indistinguishable

define a distance
in event space

eg.

$$d(e^{(1)}, e^{(2)})^2 = [\theta_1^{(1)} - \theta_1^{(2)}]^2$$

$$[\theta_2^{(1)} - \theta_2^{(2)}]^2$$



algorithm

1. pick an event with $w < 0$
(the seed)
2. add nearest event to cell C
3. if $\sum_{i \in C} w_i < 0$ goto 2
if size of $C > \text{resolution}$, goto 5

$$4. w_i \rightarrow \frac{\sum_{j \in C} w_j}{\sum_{j \in C} |w_j|} |w_i| \quad \forall i \in C$$

5. add more events in cell
& try again

Step 3: cancel things locally rather than globally (subtr. vs slicing)

Step 5: only possible if part of generator (cf. McMule)
guaranteed to work if enough points regardless of resolution

Summary

cancel things early rather than late